

Structural rigidity and thermal conductance of nanostructured carbon: nanotubes and foams

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REPORT DOCUMENTATION PAGE				Form Approved OMB No. 0704-0188	
Public reporting burden for this collection of information is estimated to average 1 hour per response, including the time for reviewing instructions, searching existing data sources, gathering and maintaining the data needed, and completing and reviewing this collection of information. Send comments regarding this burden estimate or any other aspect of this collection of information, including suggestions for reducing this burden to Department of Defense, Washington Headquarters Services, Directorate for Information Operations and Reports (0704-0188), 1215 Jefferson Davis Highway, Suite 1204, Arlington, VA 22202-4302. Respondents should be aware that notwithstanding any other provision of law, no person shall be subject to any penalty for failing to comply with a collection of information if it does not display a currently valid OMB control number. PLEASE DO NOT RETURN YOUR FORM TO THE ABOVE ADDRESS.					
1. REPORT DATE (DD-MM-YYYY) 30-05-2001		2. REPORT TYPE Workshop Presentations		3. DATES COVERED (FROM - TO) 30-05-2001 to 01-06-2001	
4. TITLE AND SUBTITLE Structural rigidity and thermal conductance of nanostructured carbon: nanotubes and foams Unclassified				5a. CONTRACT NUMBER	
				5b. GRANT NUMBER	
				5c. PROGRAM ELEMENT NUMBER	
6. AUTHOR(S) Tomanek, David ;				5d. PROJECT NUMBER	
				5e. TASK NUMBER	
				5f. WORK UNIT NUMBER	
7. PERFORMING ORGANIZATION NAME AND ADDRESS Michigan State University Physics-Astornomy Department xxxxx, MIxxxxx				8. PERFORMING ORGANIZATION REPORT NUMBER	
9. SPONSORING/MONITORING AGENCY NAME AND ADDRESS Office of Naval Research International Field Office Office of Naval Research Washington, DCxxxxx				10. SPONSOR/MONITOR'S ACRONYM(S)	
				11. SPONSOR/MONITOR'S REPORT NUMBER(S)	
12. DISTRIBUTION/AVAILABILITY STATEMENT APUBLIC RELEASE					
13. SUPPLEMENTARY NOTES See Also ADM001348, Thermal Materials Workshop 2001, held in Cambridge, UK on May 30-June 1, 2001. Additional papers can be downloaded from: http://www-mech.eng.cam.ac.uk/onr/					
14. ABSTRACT Structural rigidity, phonon modes, and thermal conductivity of carbon nanostructures: nanotubes and graphitic foam.					
15. SUBJECT TERMS					
16. SECURITY CLASSIFICATION OF:		17. LIMITATION OF ABSTRACT		18. NUMBER OF PAGES	
a. REPORT Unclassified		b. ABSTRACT Unclassified		c. THIS PAGE Unclassified	
		Public Release		28	
				19. NAME OF RESPONSIBLE PERSON Fenster, Lynn lfenster@dtic.mil	
				19b. TELEPHONE NUMBER International Area Code Area Code Telephone Number 703767-9007 DSN 427-9007	
					Standard Form 298 (Rev. 8-98) Prescribed by ANSI Std Z39.18

Acknowledgements

Savas Berber, *Michigan State University*

Young-Kyun Kwon, *U.C. Berkeley*

Susumu Saito, *Tokyo Institute of Technology*

Koichiro Umemoto, *Tokyo Institute of Technology*

Financial Support:

Office of Naval Research

DARPA

NASA

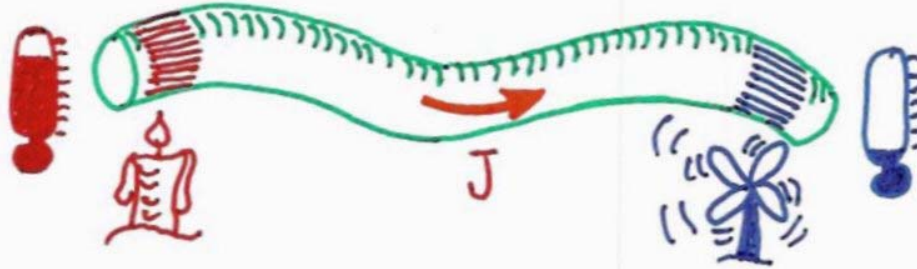


Outline

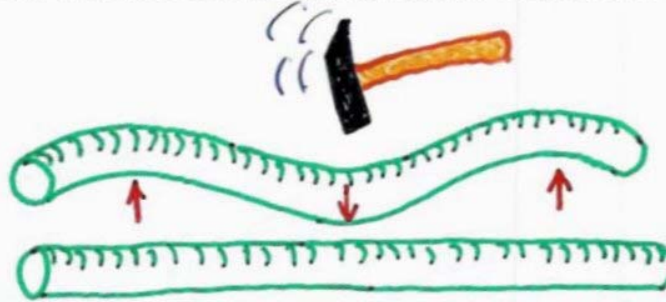
- **Motivation**
- **Computational Techniques**
- **Thermal Conductivity in Nanotubes**
 - Approaches and Pitfalls
 - Temperature Dependent Thermal Conductivity
 - Effect of Inter-Wall Interactions on Thermal Conductivity
 - Effect of Isotope Mixture on Thermal Conductivity
- **Low-Frequency Phonon Modes of Nanotubes**
 - Equilibrium Structure of Multi-Wall Tubes and Ropes
 - Soft Phonon Modes of Single-Wall Tubes
 - Libration and Rotation Dynamics of Interacting Tubes
 - Dynamics of Twistons and Orientational Melting
- **Hybrid Graphitic Foam**
 - Equilibrium Structure
 - Elastic Properties
 - Electronic Structure
- **Summary and Conclusions**

Motivation

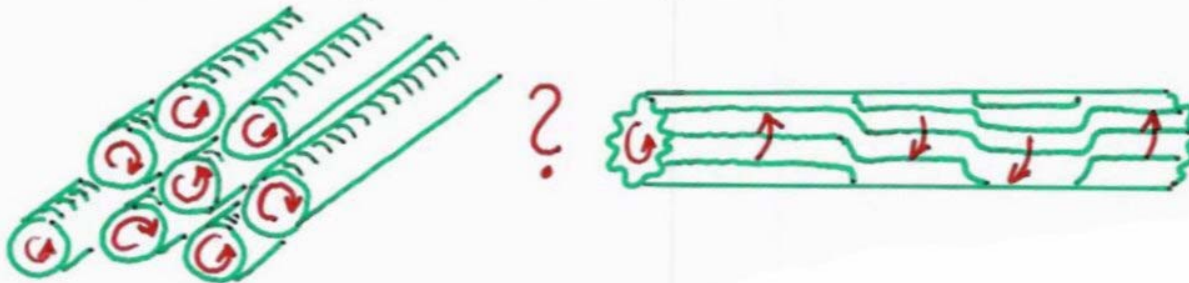
- Can Nanotubes **conduct heat** well?



- Which **Phonon Modes** limit Thermal Conduction?



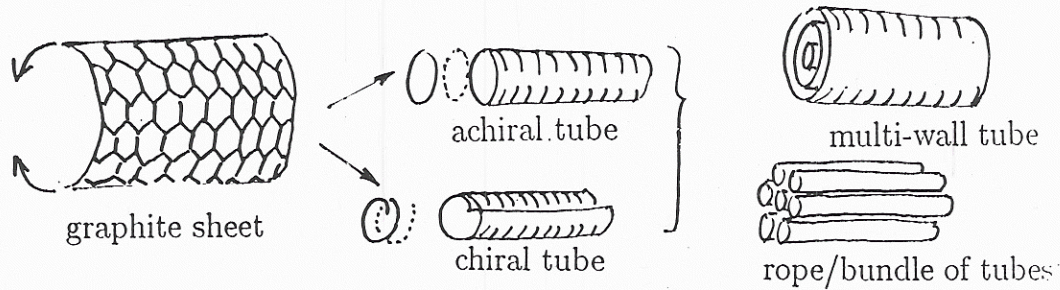
- Can interacting **Nanotubes Spin and Twist**?



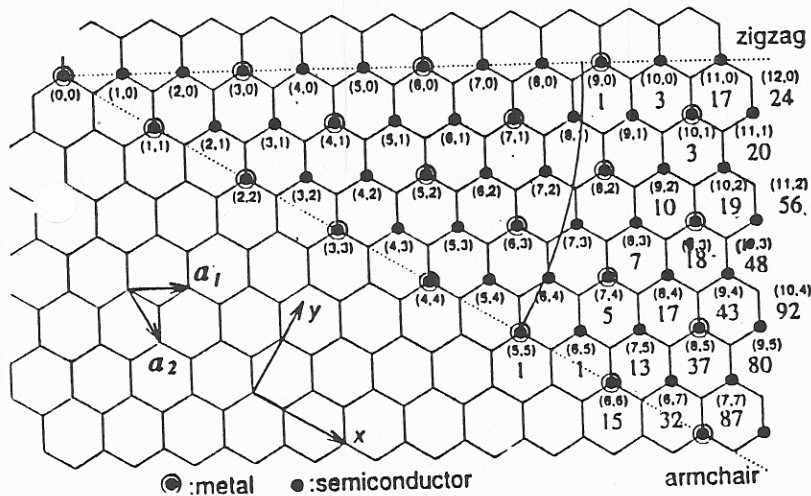
- Is nanostructured carbon foam the ultimate structural and thermal material?

■ Morphology of carbon nanotubes

• Structure of carbon nanotubes:



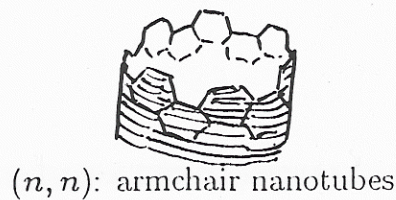
• Classification of carbon nanotubes:

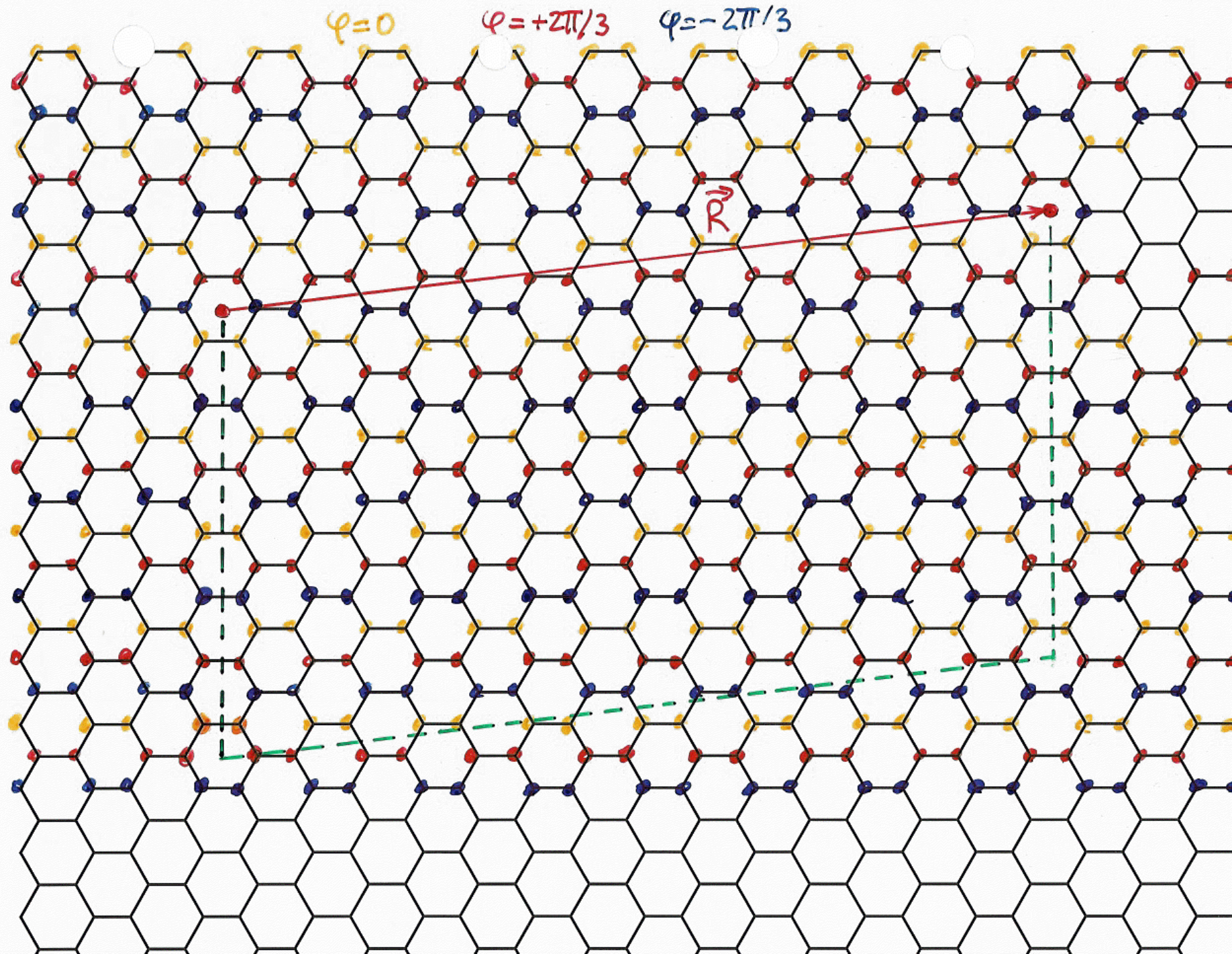


$$\vec{R} = m\vec{a}_1 + n\vec{a}_2$$

defines uniquely the (m, n) nanotube

• Types of achiral nanotubes:





Computational Techniques

- **Total Energy Calculations**

- *Ab initio* Density Functional Formalism
 - * **Plane-wave code** with energy cutoff of 50 Rydbergs
 - * Soft nonlocal pseudopotentials and Ceperley-Alder **exchange-correlation** potential
 - * 200,000 plane waves in the basis set
- Parametrized Linear Combination of Atomic Orbitals (LCAO) Formalism
 - * Parametrization based on *ab initio* Density Functional results
 - * **Very fine k-points mesh** for the inter-tube interaction
 - * $O(N)$ technique for **very large systems**
 - * Ideally suitable for **massively parallel computation**
- Tersoff Potential

- **Thermodynamics: Monte Carlo simulations**

- **Time-dependent processes:**
Molecular dynamics simulations

- **Transport: Landauer-Büttiker Formalism**

- Conductance $G(E)$ evaluated by using

$$G(E) = \frac{2e^2}{h} T(E) = \frac{1}{12.9 \text{ k}\Omega} T(E)$$

Thermal Conductivity in Nanotubes

VOLUME 84, NUMBER 20

PHYSICAL REVIEW LETTERS

15 MAY 2000

Unusually High Thermal Conductivity of Carbon Nanotubes

Savas Berber, Young-Kyun Kwon,* and David Tománek

Department of Physics and Astronomy, and Center for Fundamental Materials Research, Michigan State University,
East Lansing, Michigan 48824-1116

(Received 23 February 2000)

$$\frac{1}{A} \frac{dQ}{dt} = -\lambda \frac{dT}{dz}$$

λ = thermal conductivity

C = heat capacity/volume

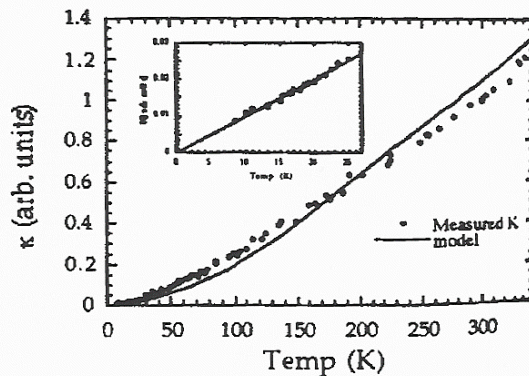
v = speed of sound

l = phonon mean free path

$$\lambda \propto C v l$$



Experimental Observation for a Nanotube “Mat”



J. Hone, M. Whitney, A. Zettl,
Synthetic Metals **103**, 2498 (1999)

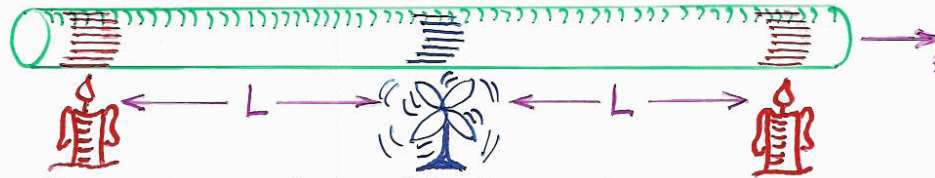
– Raw data ($T = 300$ K): $\lambda \approx 0.7$ W/m·K

– Extrapolated data ($T = 300$ K): $\lambda \approx 1,800 - 6,000$ W/m·K

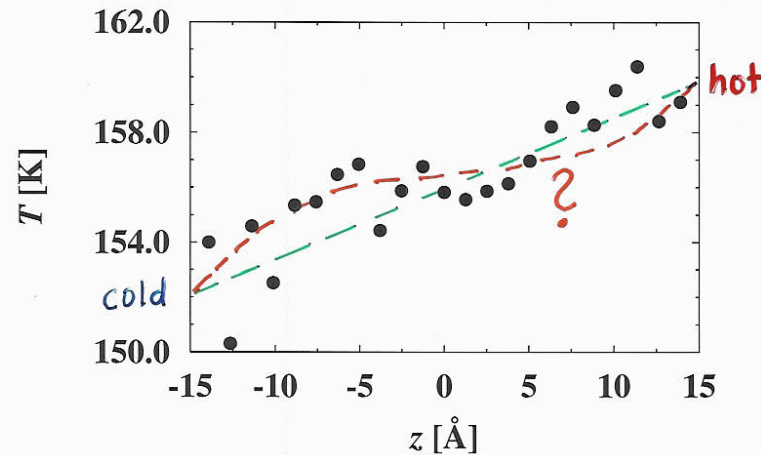
R-NT-K69

► Approaches and Pitfalls

• Direct Molecular Dynamics Approach



– Temperature profile in a (10, 10) nanotube



• Pitfalls of Direct Molecular Dynamics:

- $L \ll l \approx 100$ nm
 \Rightarrow velocity rescaling mimics “grain boundary”, limits l
- Large statistical uncertainty
- $T(z)$ not linear
- $dT/dz \gg 0 \Rightarrow$ spatial variation of thermal conductivity

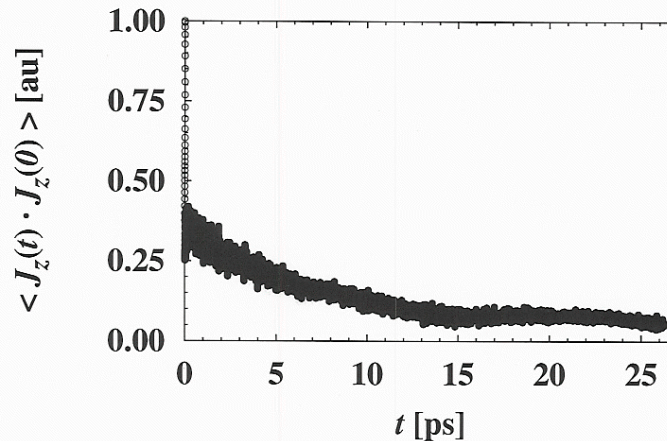
- **Equilibrium Molecular Dynamics Approach**

- Green-Kubo formula for the Navier-Stokes thermal conductivity:

$$\lambda = \frac{1}{3Vk_B T^2} \int_0^\infty \langle \mathbf{J}(t) \cdot \mathbf{J}(0) \rangle dt$$

$$\mathbf{J}(t) = \frac{d}{dt} \sum_i \mathbf{r}_i \Delta e_i = \sum_i \mathbf{v}_i \Delta e_i - \sum_i \sum_{j(\neq i)} \mathbf{r}_{ij} (\mathbf{f}_{ij} \cdot \mathbf{v}_i)$$

- Heat flux autocorrelation function in a (10, 10) nanotube



- **Pitfalls of Equilibrium Molecular Dynamics:**

- Slow convergence
- Extensive ensemble averaging crucial
- Autocorrelation function requires large data storage

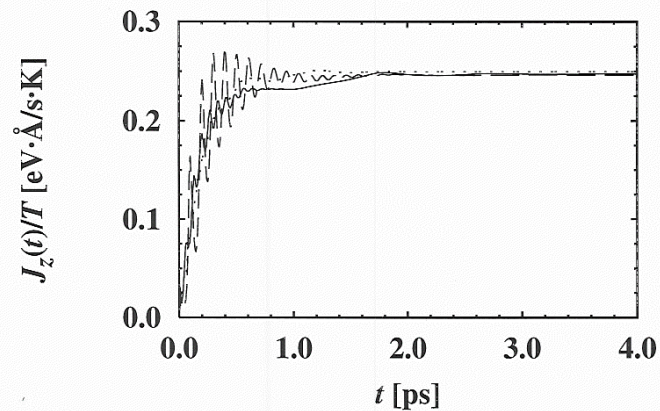
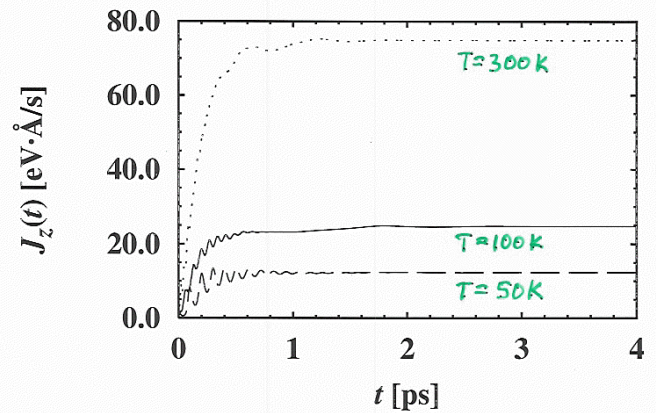
- Non-Equilibrium Molecular Dynamics Approach

$$\lambda = \lim_{F_e \rightarrow 0} \tilde{\lambda}(\mathbf{F}_e), \quad \tilde{\lambda}(\mathbf{F}_e) \equiv \lim_{t \rightarrow \infty} \frac{\langle J_z(\mathbf{F}_e, t) \rangle}{F_e T V}$$

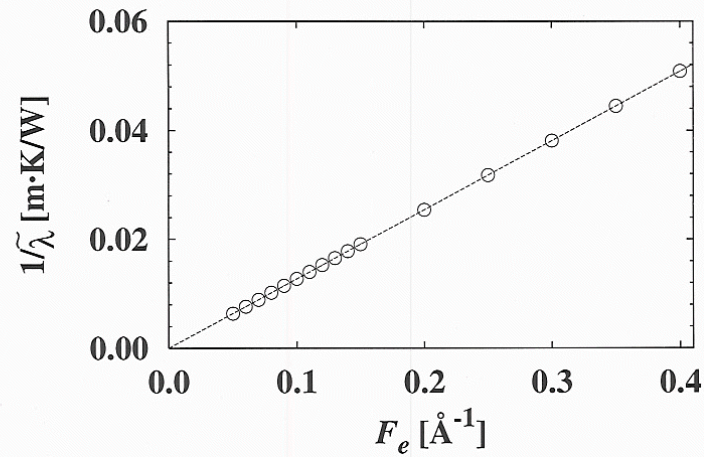
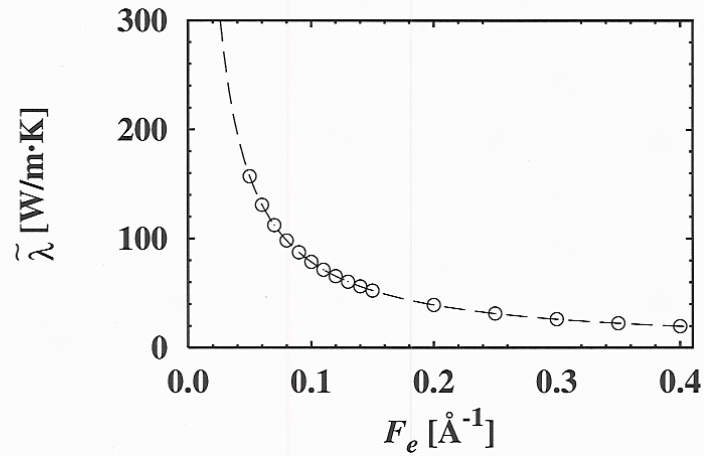
F_e = thermal force, α = Nosé-Hoover thermostat multiplier

$$\Delta \mathbf{F}_i = \Delta e_i \mathbf{F}_e - \sum_{j(\neq i)} \mathbf{f}_{ij}(\mathbf{r}_{ij} \cdot \mathbf{F}_e) + \frac{1}{N} \sum_j \sum_{k(\neq j)} \mathbf{f}_{jk}(\mathbf{r}_{jk} \cdot \mathbf{F}_e) - \alpha \mathbf{p}_i$$

– Axial heat flux in a (10,10) nanotube for $F_e = 0.2 \text{ \AA}^{-1}$



– Heat transport dependence on F_e

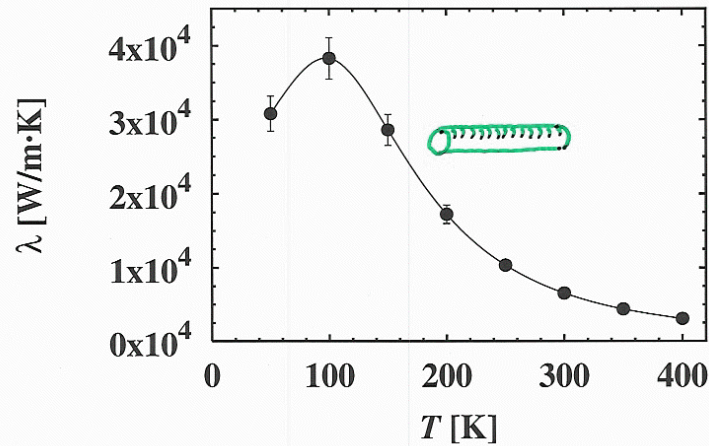


- Pitfalls of Non-Equilibrium Molecular Dynamics:

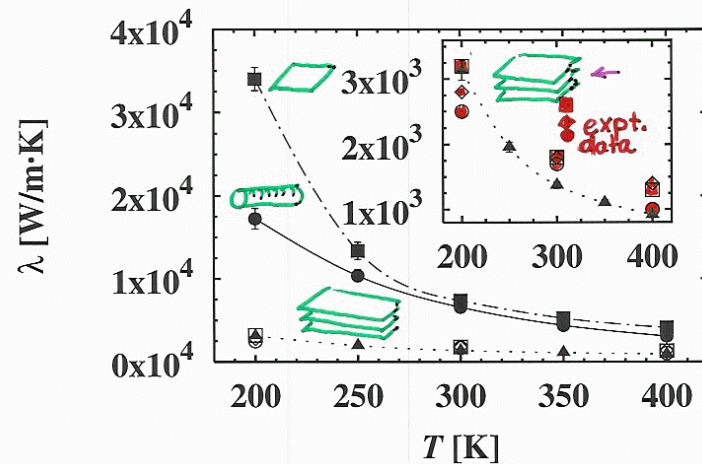
- $F_e \rightarrow 0$ limit requires careful evaluation
- Careful coupling to thermostat required

► Temperature Dependent Thermal Conductivity

- Thermal conductivity of an isolated (10,10) carbon nanotube

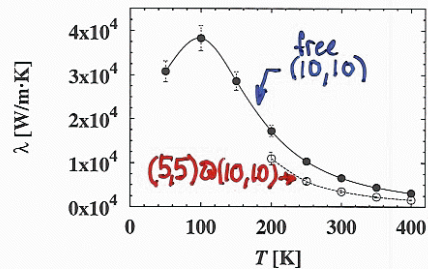


- Thermal conductivity of various carbon allotropes

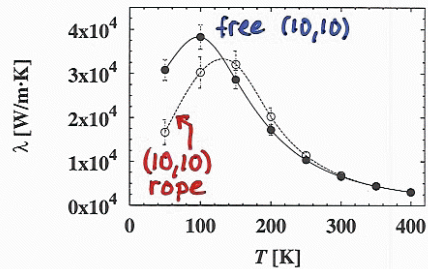


► Effect of Inter-Wall Interactions on Thermal Conductivity

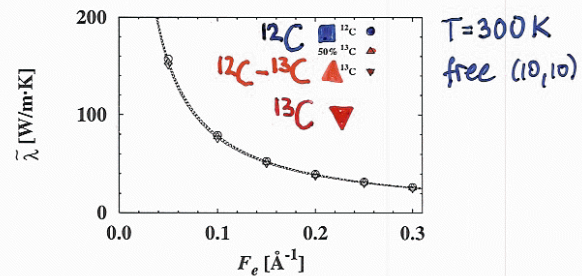
- Thermal conductivity of a (5,5)@(10,10) double-wall carbon nanotube



- Thermal conductivity of bundled (10,10) nanotubes



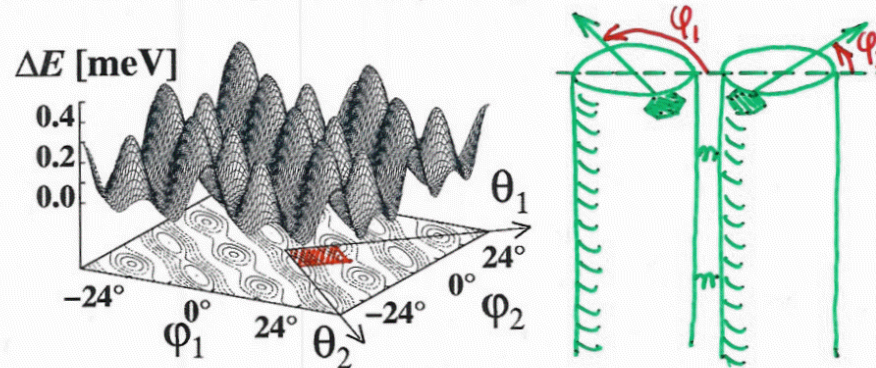
► Effect of Isotope Mixture on Thermal Conductivity



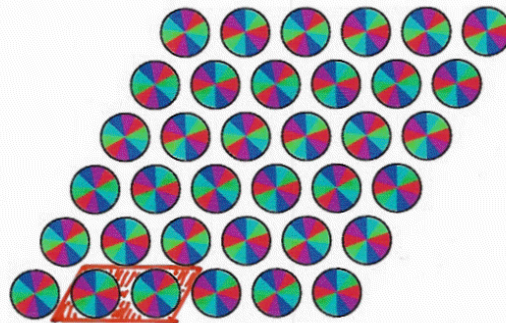
Low-Frequency Phonon Modes of Nanotubes

► Equilibrium Structure of Bundled and Multi-Wall Tubes

- Interaction between neighboring (10,10) tubes



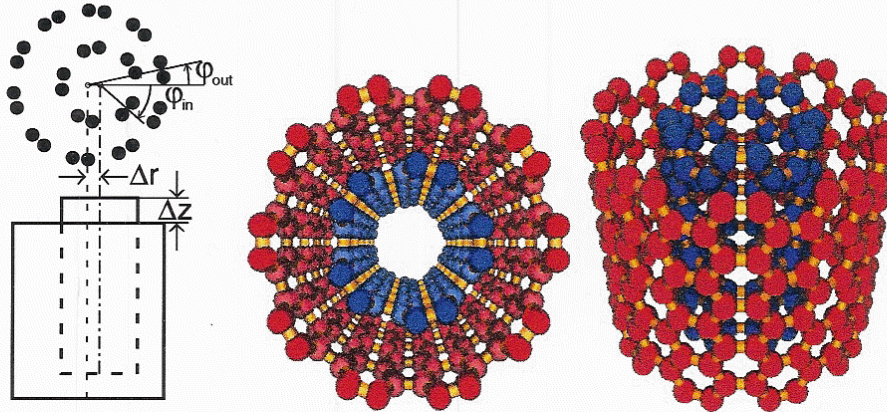
- Equilibrium structure of bundled rigid (10,10) tubes



– Two-dimensional oblique lattice.

- (5,5)@(10,10) Double-Wall Tube Geometry

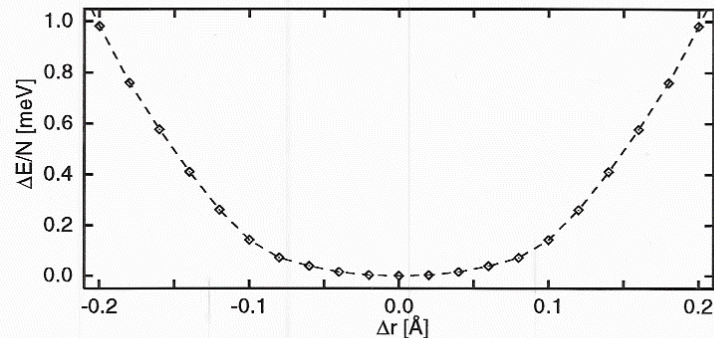
[*Electronic and structural properties of multiwall carbon nanotubes*, Young-Kyun Kwon and David Tománek, Phys. Rev. B **58**, R16001 (1998)]



(5,5)@(10,10) Double-Wall Tube

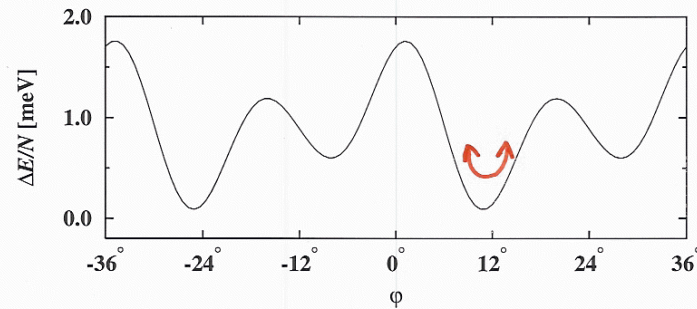
$$\begin{aligned}\Delta r &= 0 & \Delta z &= 0 \\ \varphi_{in} &= 0^\circ & \varphi_{out} &= 0^\circ\end{aligned}$$

– Dependence of Energy on Off-Axis Displacement Δr



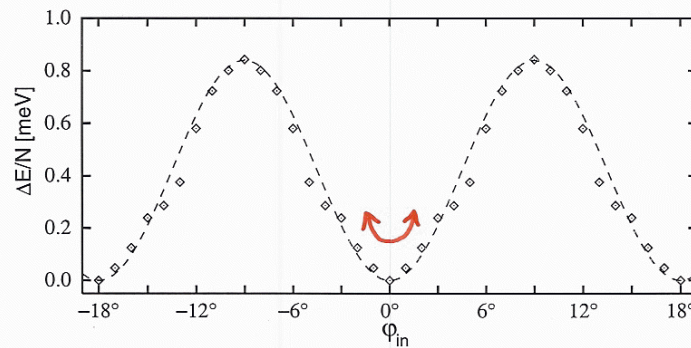
► Libration and Rotation Dynamics

- (10,10) tube “ropes” [*PRB* 58, *R13314* (1998)]



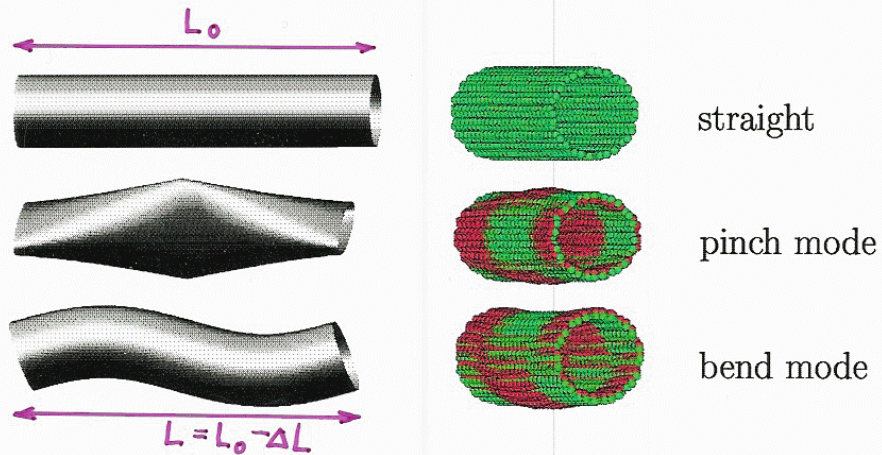
- Librational motion at $T \leq T_{OM}$
- Libration frequency $\nu \approx 12 - 60 \text{ cm}^{-1}$
- Experimental observation (W. Holmes *et al.*):
 - * Infrared modes at $\nu \approx 15, 22, 40 \text{ cm}^{-1}$
 - * These modes disappear at $T \approx 30 - 180 \text{ K}$

- (5,5)@(10,10) double-wall tube [*PRB* 58, *R16001* (1998)]

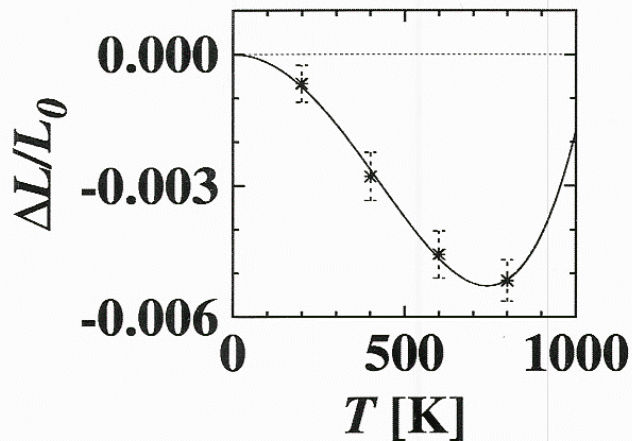


- Librational motion at $T \leq T_{OM}$
- Libration frequency $\nu \approx 33 \text{ cm}^{-1}$

► Soft Phonon Modes of Single-Wall Tubes



► Thermal Expansion of Single-Wall Tubes



⇒ Thermal contraction dominates over bond expansion at low temperatures

► Thermodynamics of Orientational Melting

VOLUME 84, NUMBER 7

PHYSICAL REVIEW LETTERS

14 FEBRUARY 2000

Orientational Melting in Carbon Nanotube Ropes

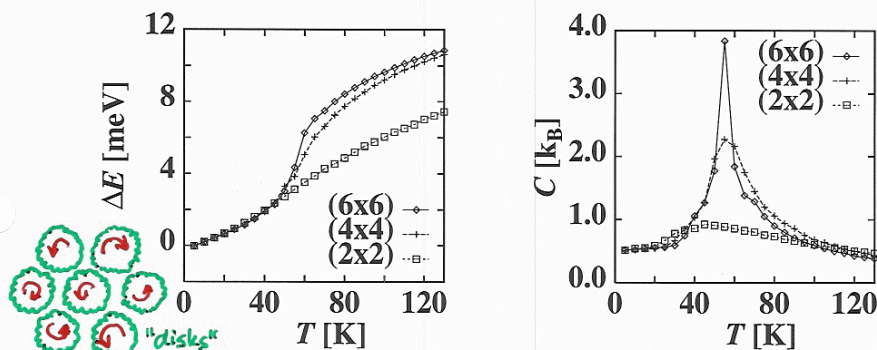
Young-Kyun Kwon and David Tománek

Department of Physics and Astronomy, and Center for Fundamental Materials Research, Michigan State University,
East Lansing, Michigan 48824-1116

(Received 24 August 1999)

• Floppy tubes – unrealistic model

(Torsionally decoupled “layers” within tube)

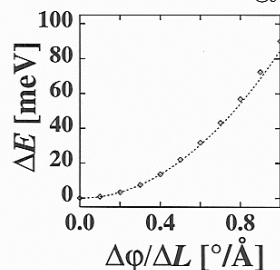


- Orientational melting at $T_{OM} \approx 55$ K
- Librational motion for $T \lesssim T_{OM}$
- Free rotation at $T \gtrsim T_{OM}$

• Rigid tubes – realistic model

(torsional coupling between “layers” within tube)

– Torsion energy



$$\Delta E_{tor,at}(\Delta \phi / \Delta L) \approx \kappa \left(\frac{\Delta \phi}{\Delta L} \right)^2$$

$$\Delta E_{tor,tot}(\Delta \phi / \Delta L) \sim \frac{1}{\Delta L} \rightarrow 0$$

for $\Delta L \rightarrow \infty$

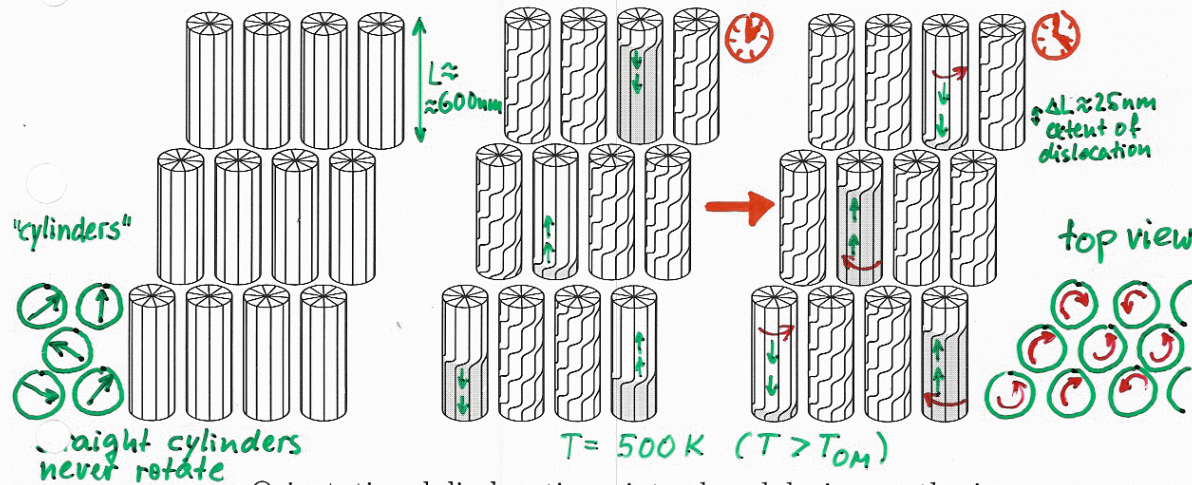
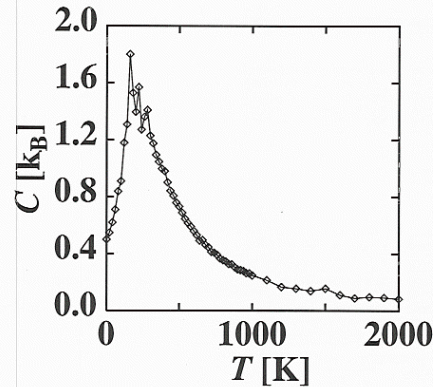
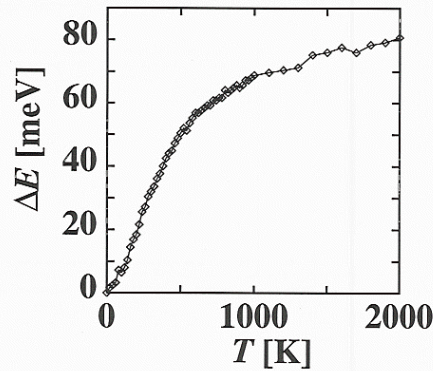


25nm orientational dislocation = twiston R-NT-168a

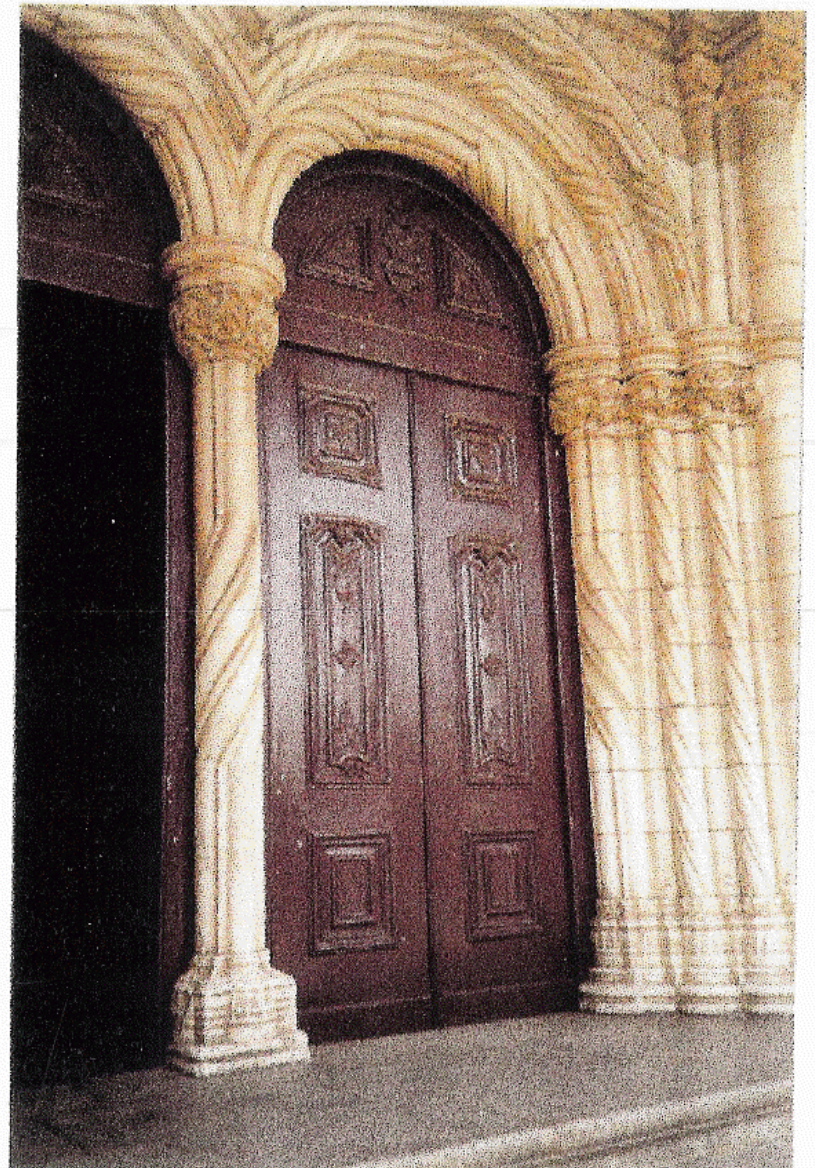
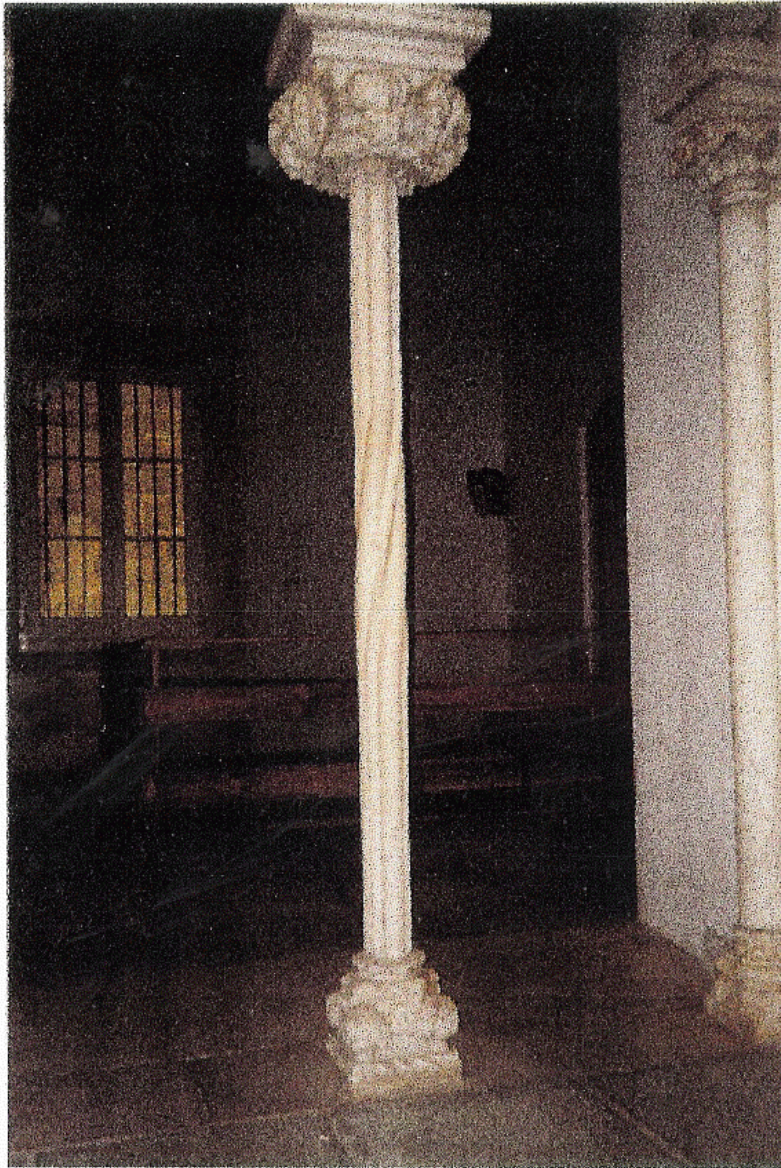


* Torsion is thermally activated during synthesis

– Orientational melting



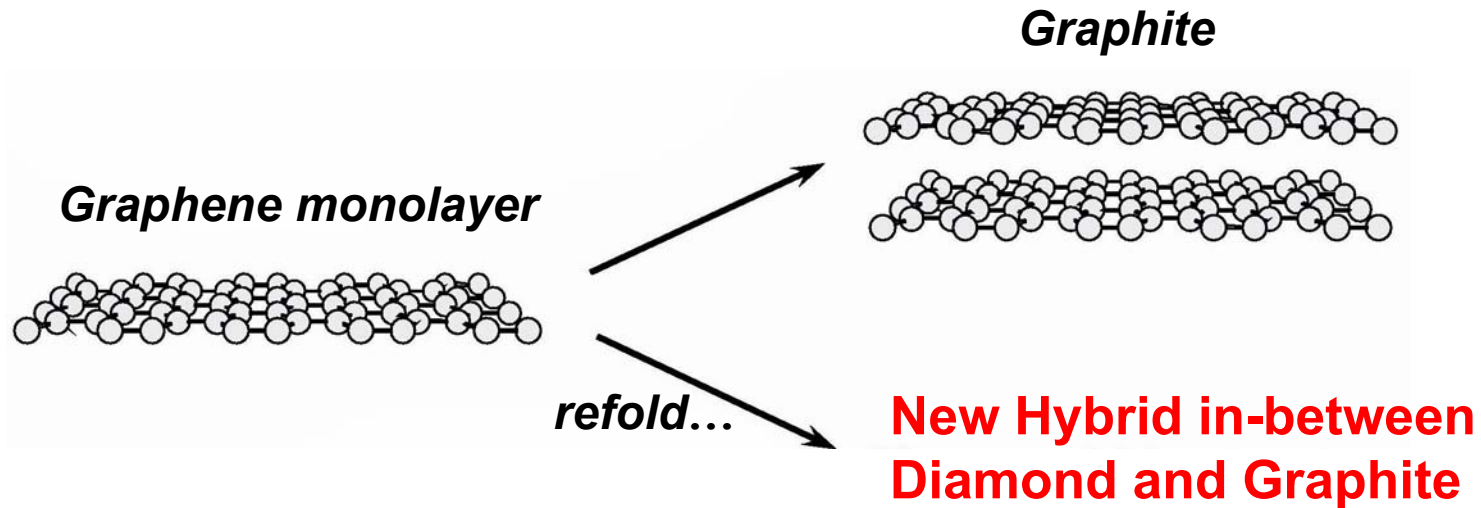
- * Orientational dislocations, introduced during synthesis, are pinned for $T \lesssim T_{OM}$.
- * Axial motion of orientational dislocations corresponds to tube rotation.
- * Orientational melting, described by free "diffusion" of orientational dislocations along tubes, occurs for $T \gtrsim T_{OM} \approx 100 \text{ K}$.
- * Orientational dislocations (twistons) are scattering centers for electrons \rightarrow strong effect on transport (1D system!).



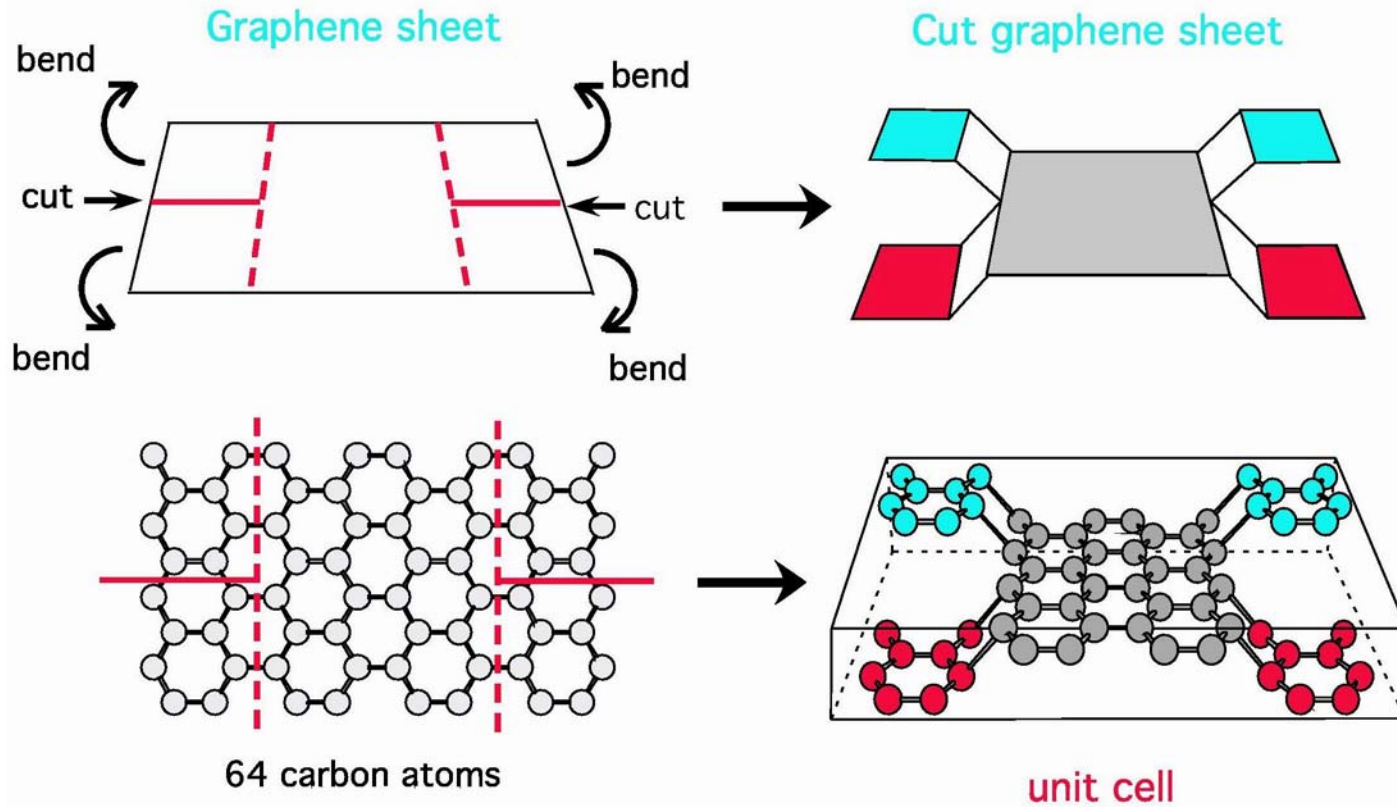
Twistons of Évora (Portugal)

Morphology of Hybrid Graphitic Foam

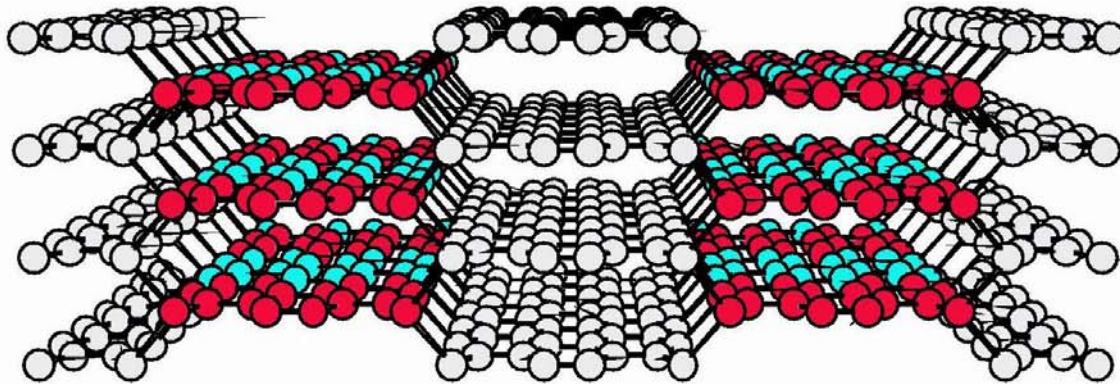
- Relationship to graphite



•Folding Process



•Hybrid Graphitic Foam

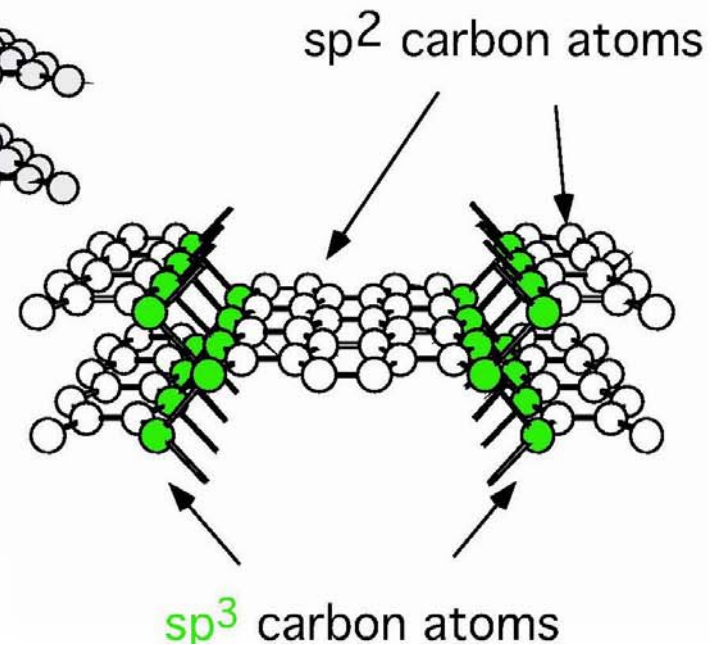


Terrace size $\rightarrow \infty$: **graphite**

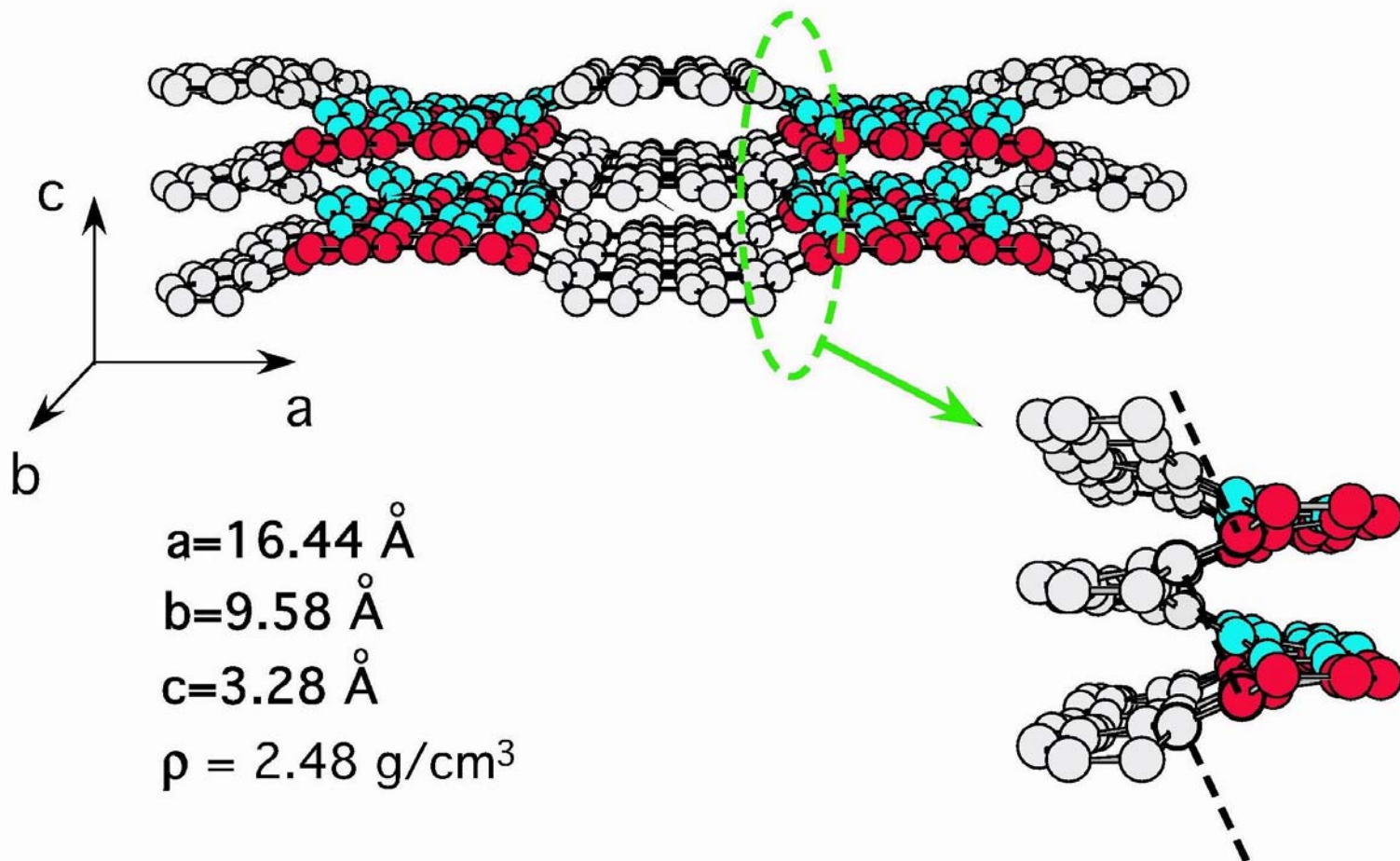
Terrace size $\rightarrow 0$: **diamond**

Present study:

- Optimized structure*
- Elastic properties*
- Electronic structure*



Optimized Structure

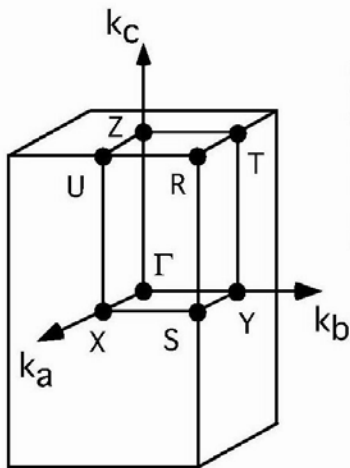


Elastic Properties

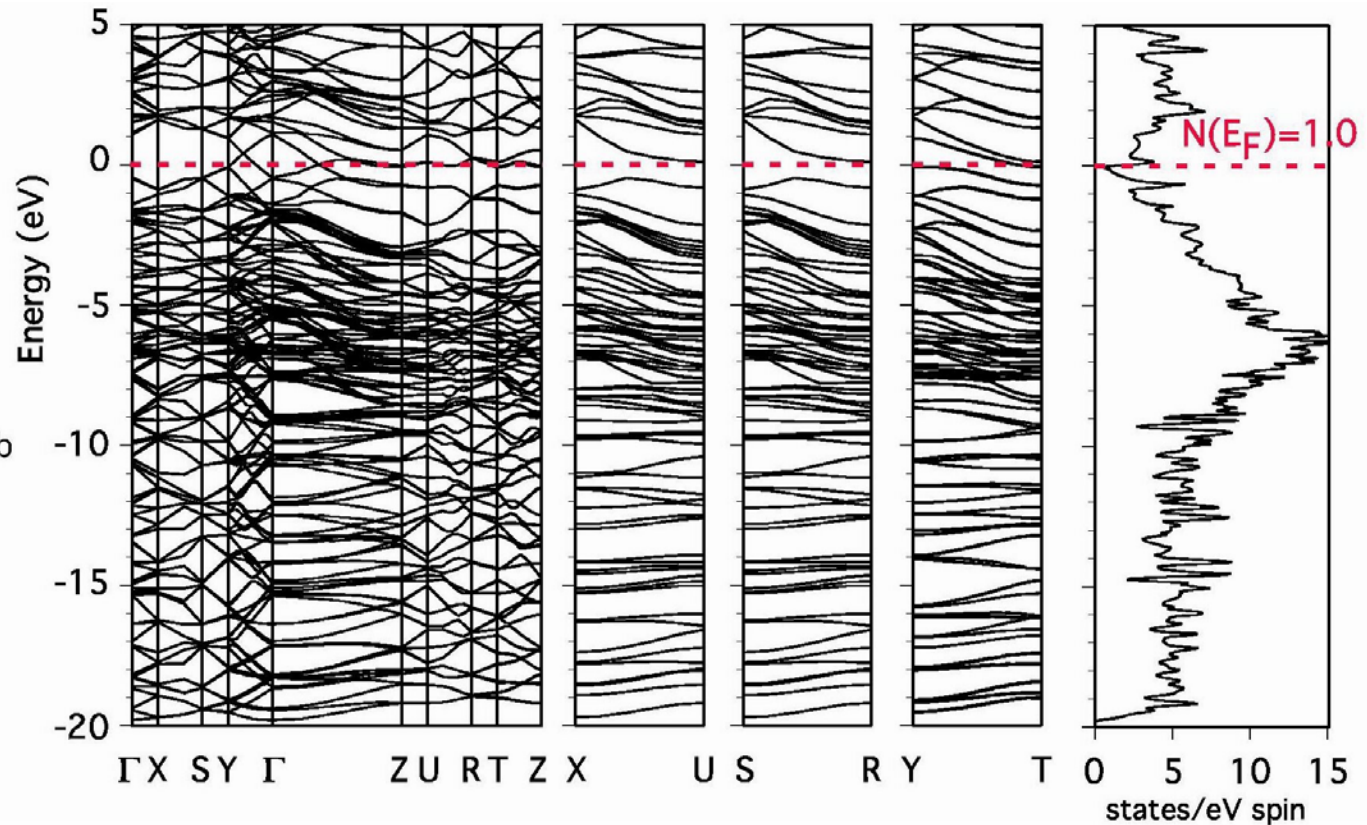
	New graphitic foam	Graphite	Diamond
ρ	2.48 g/cm ³	2.27 g/cm ³	3.51 g/cm ³
C_{11} (along a axis)	9.02 Mbar	12.3 Mbar	11.29 Mbar
C_{22} (along b axis)	9.42 Mbar		
C_{33} (along c axis)	0.816 Mbar	0.34 Mbar	
Bulk modulus	0.792 Mbar	0.326 Mbar	4.69 Mbar

Electronic Structure

Brillouin Zone



Band Structure



Graphitic foam is **metallic**: $N(E_F) = 1.0$ states/eV/spin

Summary and Conclusions

- We studied the structural rigidity, phonon modes, and thermal conductivity of **carbon nanostructures**: nanotubes and graphitic foam.
- We found that
 - The unusually **high thermal conductance** of nanotubes results from the large phonon mean free path.
 - Weak inter-tube coupling results in a very **soft librational motion**.
 - **Orientational dislocations** are frozen in during the synthesis of interacting nanotubes. Relatively free diffusion motion of orientational dislocations occurs for $T \gtrsim T_{OM} \approx 160$ K, marking the onset of **orientational melting**.
 - **Hybrid graphitic foam** offers an **unusual combination of properties**: low mass density, high structural rigidity, possibly high electric and thermal conductivity.